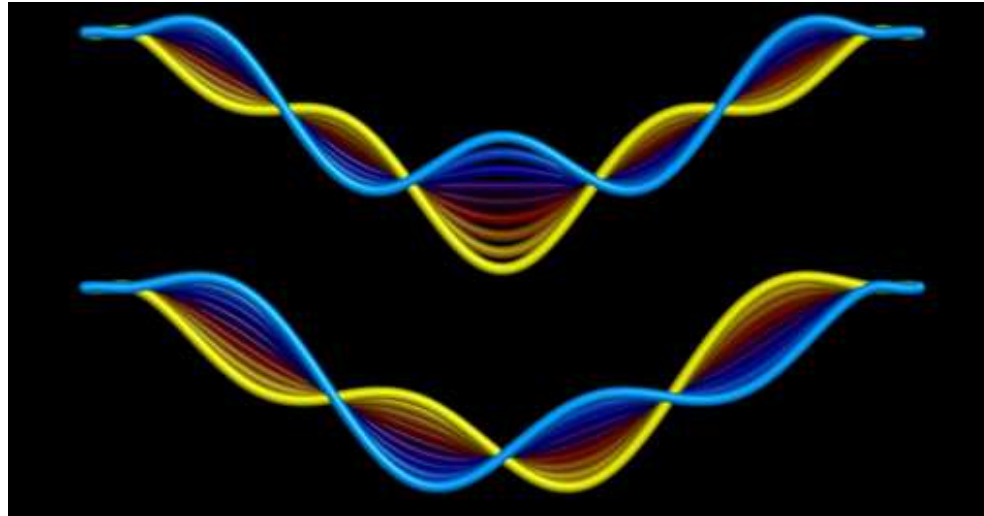


Modelling a Suspended Carbon Nanotube Oscillator

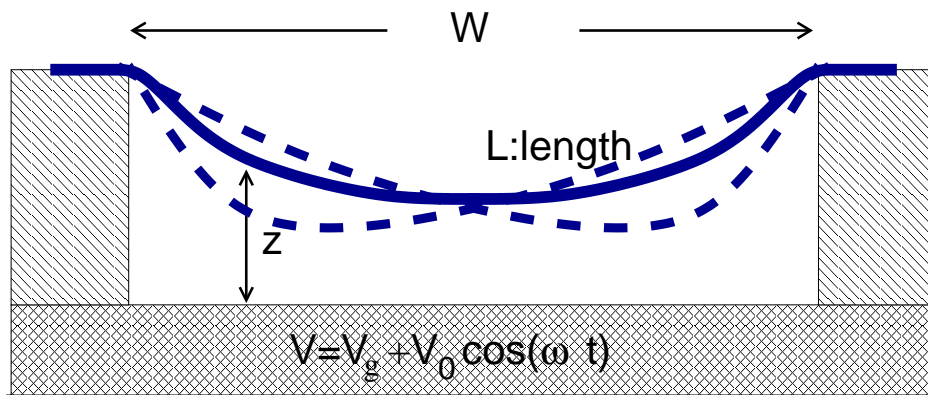
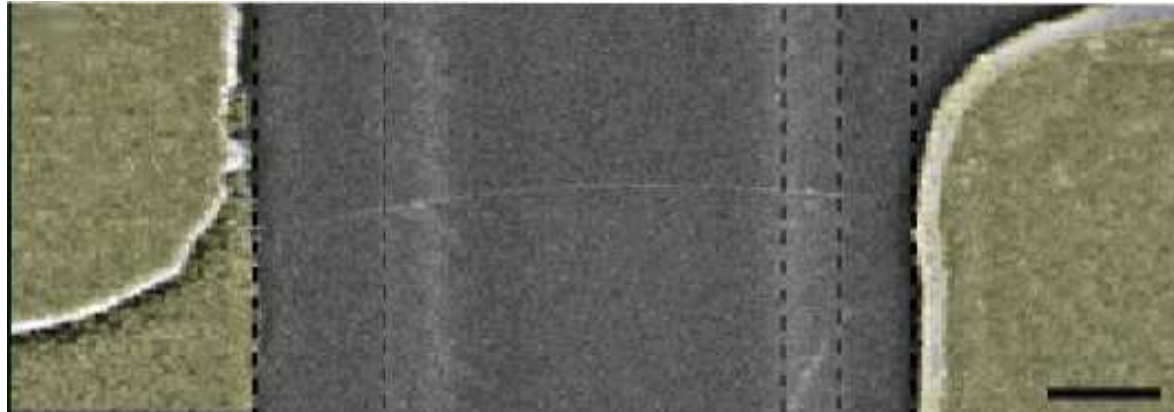


Hande Üstünel, David Roundy, Tomás Arias

Laboratory of Atomic and Solid State Physics

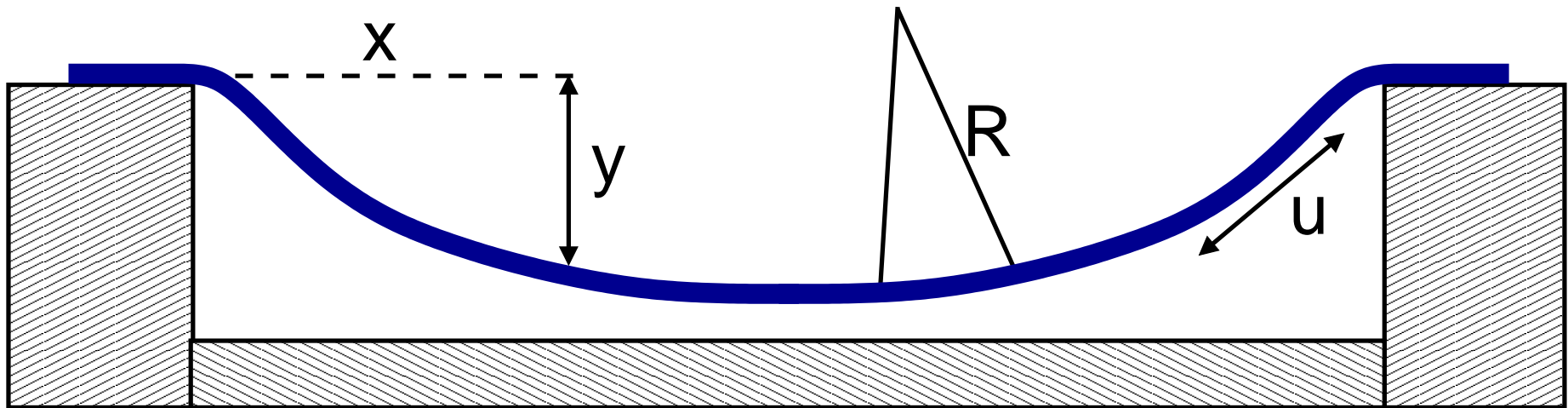
Cornell University

Nanotube suspended over trench (courtesy Vera Sazonova)



Parameters

Parameter	Typical value
Built-in slack, $s = \frac{L-W}{L}$	0–2%
Gate voltage(DC), V_g	0–6 V



- Nanotube modeled as an elastic medium with bending and stretching

$$\mathcal{U} = \frac{1}{2} \int_0^L \left[\frac{F}{R^2(x)} + E u^2(x) + f z(x) \right] dx$$

F : bending rigidity

E : extensional rigidity

f : force per length (constant)

$R(x)$: local radius of curvature

$u(x)$: local strain

$z(x)$: downward displacement

$$\mathcal{U} = \frac{1}{2} \int_0^L \left[\frac{F}{R^2(x)} + Eu^2(x) + fz(x) \right] dx$$

Small V_g

- Bending dominates.
- No extension
- ν is independent of slack, $s = \frac{L-W}{L}$.
- **Buckled beam**

Intermediate V_g

- No bending
- No extension
- $\nu \propto s^{-1/4}$
- If $s \rightarrow 0$, this limit doesn't exist.
- Hanging chain

Large V_g

- Bending is unimportant.
- Extension dominates.
- ν is independent of s .
- **Extended spring**

$$\mathcal{U} = \frac{1}{2} \int_0^L \left[\frac{F}{R^2(x)} + Eu^2(x) + fz(x) \right] dx$$

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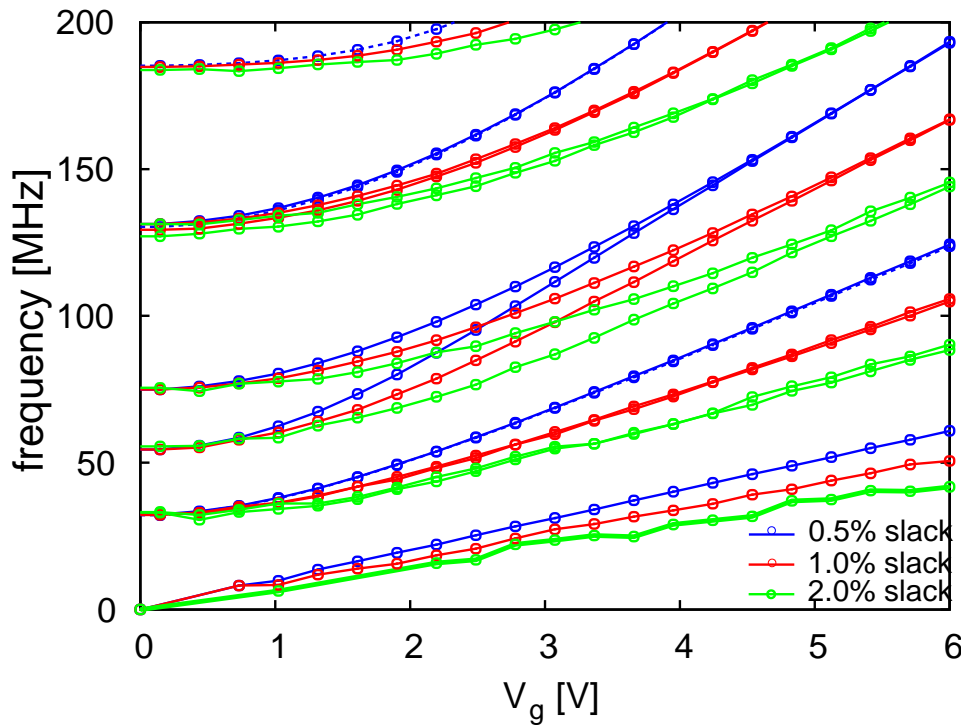
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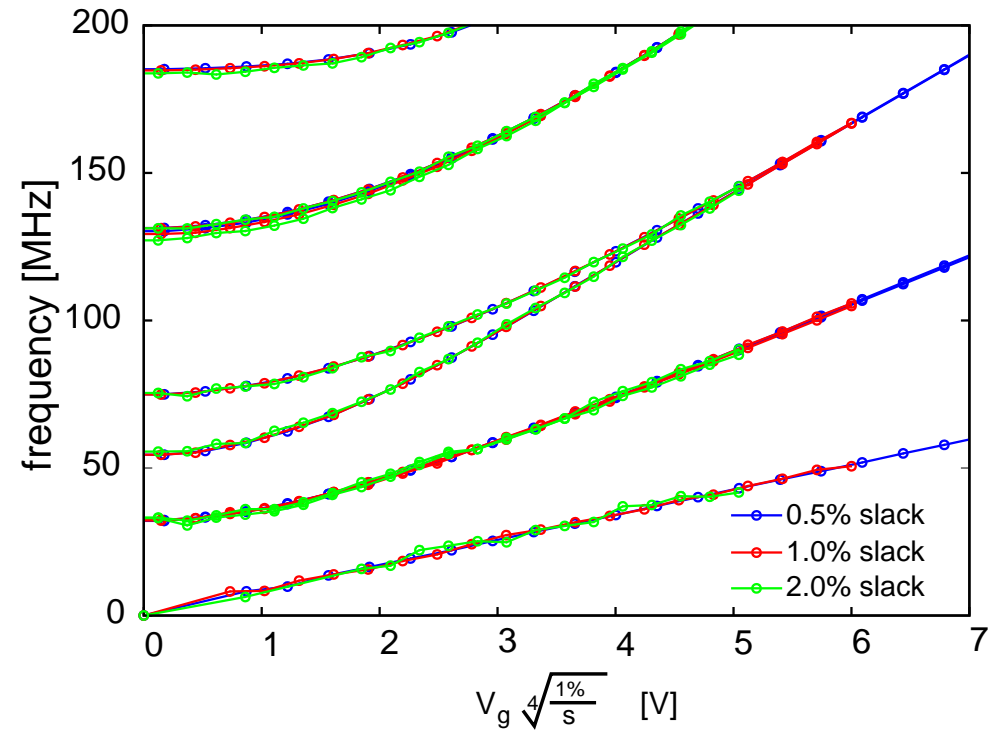
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Large V_g

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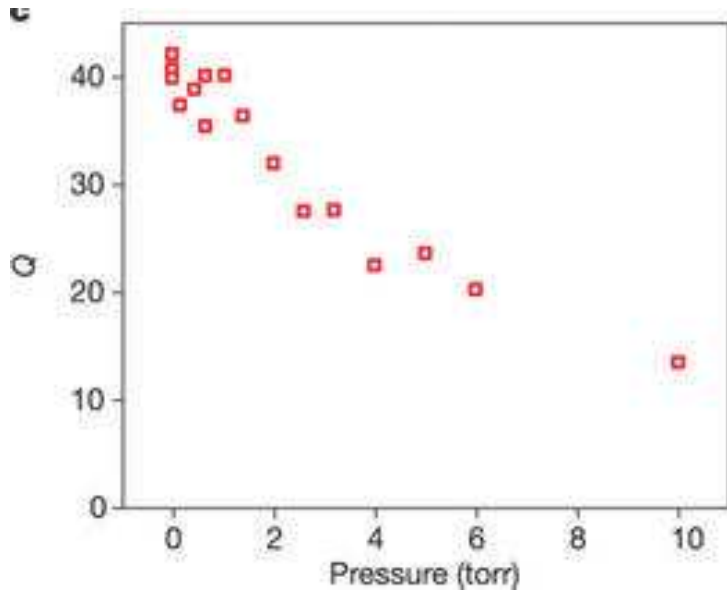
- Small V_g : ν_n independent of slack
- Large V_g : out of range



Curves collapse when scaled by $s^{-1/4}$!

- Intermediate $V_g \rightarrow \nu_n \propto s^{-1/4}$
- Small $V_g \rightarrow$ get collapsing for free

Sazonova *et al.* Nature 2004



- Phonon–phonon interactions are an intrinsic cause of mechanical loss.
- There are two limits of interest.

Nanotubes have low Q -factor

Largest observed ever $\sim 1000!$

- Acoustic phonon disturbs thermal phonon equilibrium.
- Phonons redistribute causing mechanical loss.
- Loss is characterized by inverse quality factor.

$$Q^{-1} \propto \frac{\nu\Omega}{\nu^2 + \Omega^2} \int_0^L dx \rho_{\text{slack}}^2(x) \rho_{\text{mode}}^2(x) \int \frac{dq}{2\pi} (\gamma_q n_q^{eq})^2$$

- γ_q describes change in phonon frequencies with radius of curvature.

$$\omega_q = \omega_q^0 + \gamma_q/R^2$$

Ω : frequency of acoustic mode ν : relaxation rate of thermal phonons

$\rho_{\text{slack}}(x)$: curvature of slack profile $\rho_{\text{mode}}(x)$: curvature of acoustic mode

q : wavevector of thermal phonons n_q^{eq} : equilibrium distribution of thermal phonons

- Ballistic collisions between thermal phonons and acoustic phonon cause energy loss.
- Slack and acoustic mode can be treated as athermally populated phonon modes.
- Momentum conservation is satisfied only at band crossings.
- Inverse quality factor :

$$Q^{-1} \propto n_q(n_q + 1) \frac{1}{\Delta v_g} \frac{1}{\omega_q^2 \Omega^2} \underbrace{\left(\frac{\partial \mathcal{U}^4}{\partial^2 \rho \partial^2 \eta_q} \right)^2}_K$$

- K causes band gap opening at band crossings.

$$\Delta\omega_q = K / m\omega R^2$$

Ω : frequency of acoustic mode

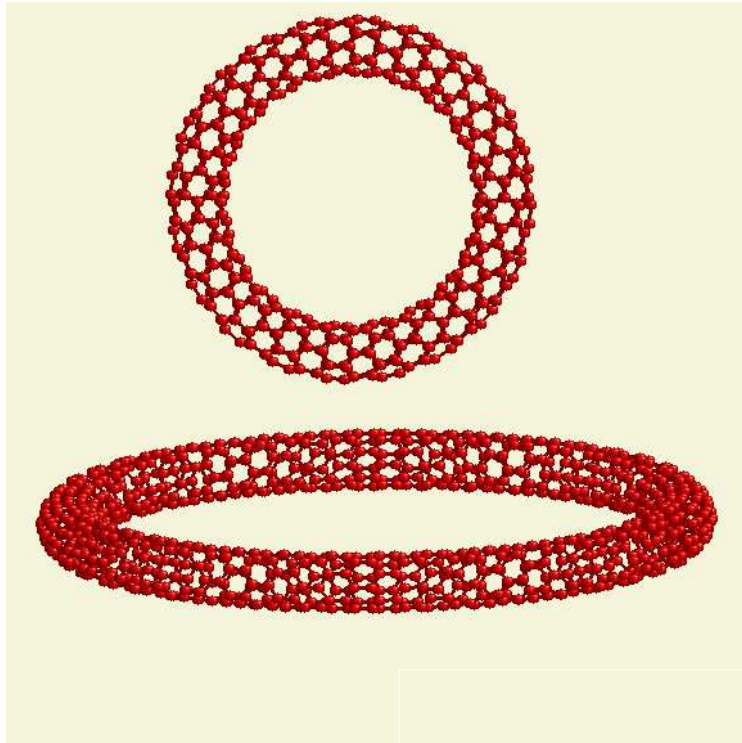
ω_q : frequency of thermal phonons

ρ : radius of curvature

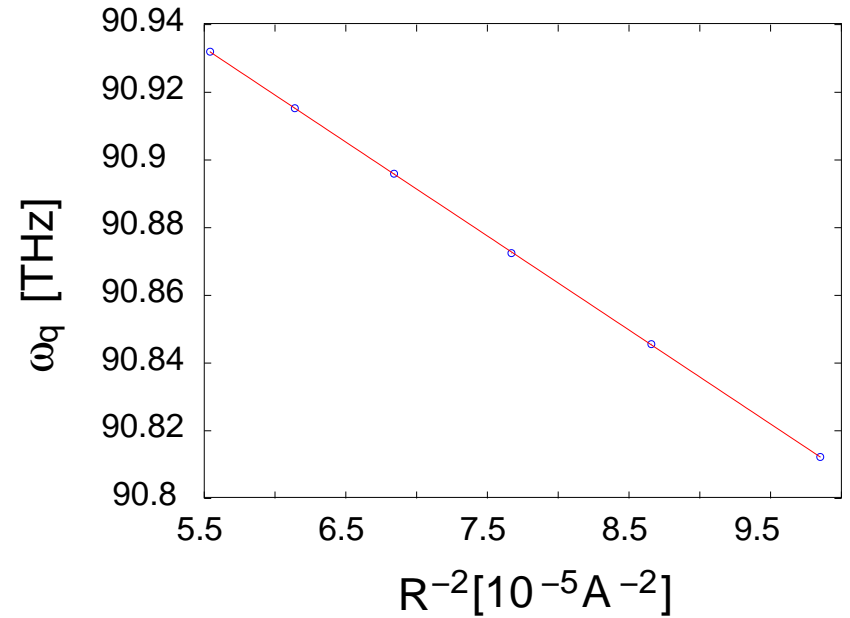
Δv_g : group velocity difference of thermal phonons

η_q : amplitude in thermal phonon modes

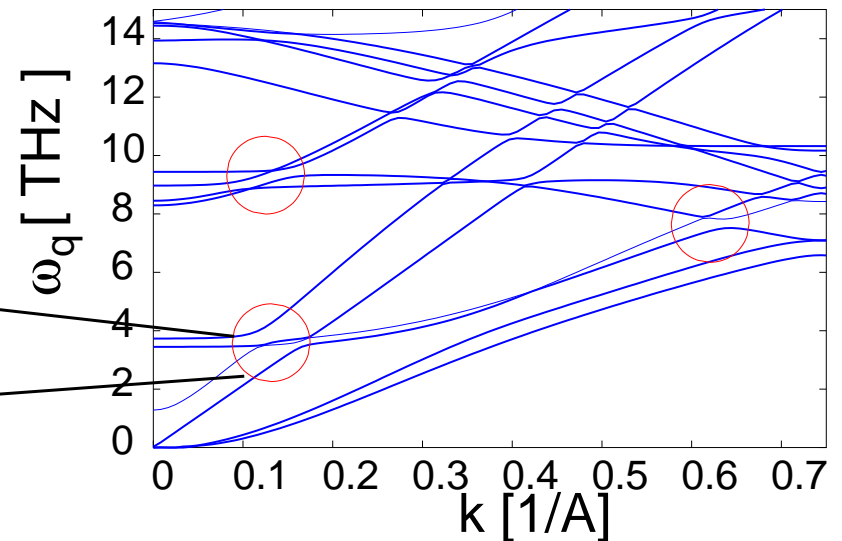
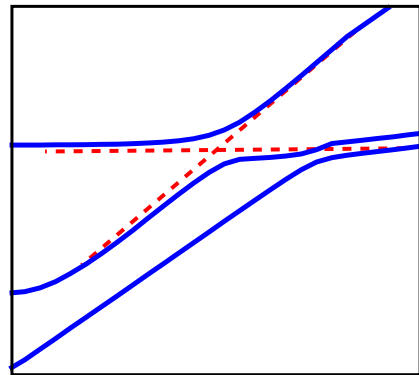
n_q^{eq} : equilibrium distribution of thermal phonons



- γ_q is the slope of ω_q vs R^{-2} plot.



- K causes band gap opening.



Regime	Quality factor (room temperature)	Quality factor (4 K)
Akhieser	$\approx 10^{14}$	$\approx 10^{63}$
Landau-Rümer	$\approx 10^5$	$\approx 10^{23}$

- We have used a continuum model to study suspended nanotubes.
- We have studied mechanical loss due to phonon-phonon coupling.
- Our results indicate that this intrinsic loss mechanism can be ruled out, especially at low temperatures.
- Other possible sources of loss include
 - Clamping
 - Interaction with gas molecules
 - Residual fabrication material left on the nanotube
 - Electron-phonon interaction